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## LETTER TO THE EDITOR

# Markov processes and a multiple generating function of product of generalized Laguerre polynomials 

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#### Abstract

From the spectral representation of the transition probability of birth-and-death processes, Karlin and McGregor show that the transition probability for the infinite server Markovian queue is in the form of a diagonal sum involving a product of Charlier polynomials. By using Meixner's bilinear generating formula for the Charlier polynomials and the Markov property, a multiple generating for the Charlier polynomials is deduced from the ChapmanKolmogorov equation. The resulting formula possesses the same genre of a multiple generating function for the generalized Laguerre polynomials discussed by Messina and Paladimo, the explicit solution of which is recently given by the present author.


## 1. Introduction

In a previous communication [1], we discussed the explicit evaluation of a multiple generating function involving the product of generalized Laguerre polynomials

$$
L_{n}^{(\alpha)}(x)=\frac{(\alpha+1)_{n}}{n!}{ }_{1} F_{1}(-n ; \alpha+1 ; x)
$$

of the form

$$
\begin{equation*}
I_{n}=\sum_{s_{1}, s_{2}, \ldots, s_{n}=0}^{\infty} z_{1}^{s_{1}} z_{2}^{s_{2}} \cdots z_{n}^{s_{n}} L_{s_{2}}^{\left(s_{1}-s_{2}\right)}(x) L_{s_{3}}^{\left(s_{2}-s_{3}\right)}(x) \cdots L_{s_{1}}^{\left(s_{n}-s_{1}\right)}(x) \tag{1}
\end{equation*}
$$

It can be shown that $I_{n}$ is a symmetric function of the variables $z_{1}, z_{2}, \ldots, z_{n}$, and is explicitly given by

$$
\begin{equation*}
I_{n}=\frac{1}{1-\sigma_{n}} \exp -\left\{\frac{x}{1-\sigma_{n}}\left[(n-1) \sigma_{n}+\psi(1)-1\right]\right\} \quad n \geqslant 1 \tag{2}
\end{equation*}
$$

where the elementary symmetric functions are defined as

$$
\begin{aligned}
\sigma_{0} & =1 \\
\sigma_{1} & =\sum_{i=1}^{n} z_{i} \\
\sigma_{2} & =\sum_{i<j}^{n} z_{i} z_{j} \\
\sigma_{3} & =\sum_{i<j<k}^{n} z_{i} z_{j} z_{k}
\end{aligned}
$$

$$
\sigma_{n}=z_{1} z_{2} \cdots z_{n}
$$

and $\psi(t)$, the generating function of $\sigma_{i}, 0 \leqslant i \leqslant n$, is given by

$$
\psi(t)=\sum_{i=0}^{n} t^{i} \sigma_{i}=\prod_{i=1}^{n}\left(1+z_{i} t\right) .
$$

The multiple generating function in (1) was originally discussed by Messina and Paladimo [2], who gave a solution, in implicit form, by appealing to properties of the trace of a certain sequence of operators. Apart from its relevance in the theory of special functions, this generating function is also encountered in a variety of situations involving researches in physics and chemistry. (See [2] for further references.)

We will now show that an allied sum of the form of (1) is intimately connected with the transition probability of a certain stochastic process arising from the infinite server Poisson queueing system. The Markovian nature of the process, and in particular, the ChapmanKolmogorov equation, allows us to deduce an explicit expression of a sum which bears a close resemblance to the solution given in (2). The method makes use of the result of Karlin and McGregor [3, 4] on the spectral representation of the transition probability of the birth-and-death process in terms of orthogonal polynomials, the outline of which is briefly described below.

## 2. The birth-and-death process

A birth-and-death process is a stationary Markov process $\{X(t), t \geqslant 0\}$ whose state space $S$ is the set of non-negative integers $\{0,1,2, \ldots\}$, and whose transition probabilities

$$
P_{i j}(t)=\operatorname{Pr}[X(t)=j \mid X(0)=i] \quad i, j \in S \quad t \geqslant 0
$$

satisfy the following infinitesimal transition scheme as $\Delta t \rightarrow 0$

$$
P_{i j}(\Delta t)= \begin{cases}\lambda_{i} \Delta t+\mathrm{o}(\Delta t) & \text { if } j=i+1 \\ \mu_{i} \Delta t+\mathrm{o}(\Delta t) & \text { if } j=i-1 \\ 1-\left(\lambda_{i}+\mu_{i}\right) \Delta t+\mathrm{o}(\Delta t) & \text { if } j=i \\ \mathrm{o}(\Delta t) & \text { if }|j-i|>1\end{cases}
$$

where $\lambda_{i}>0$ for $i \geqslant 0, \mu_{i}>0$ for $i \geqslant 1$ and $\mu_{0} \geqslant 0$.
The infinitesimal matrix of the birth-and-death process is of the form

$$
\mathbf{A}=\left(\begin{array}{cccc}
-\left(\lambda_{0}+\mu_{0}\right) & \lambda_{0} & 0 & \cdots \\
\mu_{1} & -\left(\lambda_{1}+\mu_{1}\right) & \lambda_{1} & \cdots \\
0 & \mu_{2} & -\left(\lambda_{2}+\mu_{2}\right) & \cdots \\
\vdots & \vdots & \vdots & \vdots
\end{array}\right) .
$$

This matrix determines a system of orthogonal polynomials $\left\{Q_{n}(x)\right\}_{n=0}^{\infty}$ by means of the following three-term recurrence relation

$$
\begin{equation*}
-x Q_{n}(x)=\mu_{n} Q_{n-1}(x)-\left(\lambda_{n}+\mu_{n}\right) Q_{n}(x)+\lambda_{n} Q_{n+1}(x) \quad n \geqslant 0 \tag{3}
\end{equation*}
$$

where $Q_{n}(x)=0$ for $n<0$, and $Q_{0}(x)=1$.

It is shown in [3] that there is a positive measure $\phi(x)$ on $0 \leqslant x<\infty$ for which the following orthogonality conditions hold

$$
\int_{0}^{\infty} Q_{i}(x) Q_{j}(x) \mathrm{d} \phi(x)=\frac{\delta_{i j}}{\pi_{j}} \quad i, j \in S
$$

where $\pi_{0}=1, \quad \pi_{n}=\lambda_{0} \lambda_{1} \cdots \lambda_{n-1} / \mu_{1} \mu_{2} \cdots \mu_{n}$.
Moreover, the transition probability matrix $\mathbf{P}(t)=\left(P_{i j}(t)\right)$ of the process which is uniquely determined by $\mathbf{A}$, has the following explicit spectral representation

$$
\begin{equation*}
P_{i j}(t)=\pi_{j} \int_{0}^{\infty} \mathrm{e}^{-x t} Q_{i}(x) Q_{j}(x) \mathrm{d} \phi(x) \quad i, j \in S \tag{4}
\end{equation*}
$$

In addition, from the Markovian nature of the process, $\mathbf{P}(t)$ will satisfy the ChapmanKolmogorov equation

$$
\mathbf{P}(s+t)=\mathbf{P}(s) \mathbf{P}(t) \quad s, t \geqslant 0
$$

or equivalently

$$
\begin{equation*}
P_{i j}(s+t)=\sum_{k=0}^{\infty} P_{i k}(s) P_{k j}(t) \quad i, j \in S \tag{5}
\end{equation*}
$$

## 3. The $M / M / \infty$ queueing process

For the infinite server queue $\mathrm{M} / \mathrm{M} / \infty$ with Poisson input and exponential service, we have

$$
\begin{array}{lc}
\lambda_{n}=\lambda & n \geqslant 0 \\
\mu_{n}=n \mu & n \geqslant 1 .
\end{array}
$$

From (3), the orthogonal polynomials associated with this process, denoted by $Q_{n}(x) \equiv$ $Q_{n}(x, \lambda, \mu)$, satisfy the following three-term recurrence relation:

$$
-x Q_{n}(x)=n \mu Q_{n-1}(x)-(\lambda+n \mu) Q_{n}(x)+\lambda Q_{n+1}(x) \quad n \geqslant 0
$$

where $Q_{n}(x)=0$ for $n<0$, and $Q_{0}(x)=1$.
As shown in [5], this is identical with the Charlier polynomials $c_{n}(x ; a)$ which satisfy the following three-term recurrence relation [5]:

$$
-x c_{n}(x ; a)=n c_{n-1}(x ; a)-(n+a) c_{n}(x ; a)+a c_{n+1}(x ; a) \quad n \geqslant 0
$$

where $c_{0}(x ; a)=1, c_{n}(x ; a)=0$ for $n<0$. Thus on identification we have

$$
Q_{n}(x)=c_{n}\left(\frac{x}{\mu} ; \frac{\lambda}{\mu}\right)
$$

It is well known [6] that the Charlier polynomials are a system of discrete orthogonal polynomials with respect to the following Poisson probability distribution:

$$
p(x)=\mathrm{e}^{-a} \frac{a^{x}}{x!} \quad x=0,1,2, \ldots
$$

Hence, the spectral measure $\phi(x)$ consists of masses

$$
\mathrm{d} \phi(x)=\mathrm{e}^{-a} \frac{a^{n}}{n!} \quad \text { at } x_{n}=n \mu \quad n=0,1,2, \ldots
$$

where $a=\lambda / \mu$.
The symmetry property of the Charlier polynomials [6]

$$
c_{m}(n ; a)=c_{n}(m ; a)
$$

implies that

$$
Q_{m}(n \mu)=Q_{n}(m \mu) .
$$

From (4), we therefore have the following explicit spectral representation of the transition probability:

$$
\begin{align*}
P_{i j}(t) & =\frac{a^{j}}{j!} \sum_{r=0}^{\infty} \mathrm{e}^{-r \mu t} Q_{i}(r \mu) Q_{j}(r \mu) \mathrm{e}^{-a} \frac{a^{r}}{r!} \\
& =p(j) \sum_{r=0}^{\infty} \frac{\left(a \mathrm{e}^{-\mu t}\right)^{r}}{r!} c_{i}(r ; a) c_{j}(r ; a) \\
& =\frac{\mathrm{e}^{a(1-z)}}{j!}[a(1-z)]^{j}(1-z)^{i} c_{i}\left(j ;-\frac{a(1-z)^{2}}{z}\right) \tag{6}
\end{align*}
$$

where in arriving at (6) we have set $\mathrm{e}^{-\mu t}=z$ and made use of the following Meixner's bilinear generating formula for the Charlier polynomials [7]:

$$
\sum_{k=0}^{\infty} \frac{z^{k}}{k!} c_{k}(m ; x) c_{k}(n ; y)=\mathrm{e}^{z}\left(1-\frac{z}{x}\right)^{m}\left(1-\frac{z}{y}\right)^{n} c_{m}\left(n ;-\frac{(x-z)(y-z)}{z}\right) .
$$

## 4. A multiple generating function

For the process $\{X(t), t \geqslant 0\}$, let us define

$$
P_{s_{1} s_{1}}\left(t_{1}\right)=\operatorname{Pr}\left[X\left(t_{1}\right)=s_{1} \mid X(0)=s_{1}\right]
$$

as the one-step recurrence probability at state $s_{1}$, and similarly

$$
P_{s_{1} s_{1}}\left(\tau_{n}\right)=\operatorname{Pr}\left[X\left(\tau_{n}\right)=s_{1} \mid X(0)=s_{1}\right] \quad \tau_{n}=\sum_{i=1}^{n} t_{i}
$$

as the $n$-step recurrence probability. The sum

$$
\begin{equation*}
J_{n}=\sum_{s_{1}=0}^{\infty} P_{s_{1} s_{1}}\left(\tau_{n}\right) \tag{7}
\end{equation*}
$$

may thus be interpreted as the $n$-step recurrence probability for all the states in the time interval $\tau_{n}$.

By virtue of the Markovian nature of the process, a repeated use of the ChapmanKolmogorov equation in (5), for summations running from $s_{2}$ to $s_{n}$, also yields the following alternative expression for $J_{n}$ :

$$
\begin{equation*}
J_{n}=\sum_{s_{1}, s_{2}, \ldots, s_{n}=0}^{\infty} P_{s_{1} s_{2}}\left(t_{1}\right) P_{s_{2} s_{3}}\left(t_{2}\right) \cdots P_{s_{n-1} s_{n}}\left(t_{n-1}\right) P_{s_{n} s_{1}}\left(t_{n}\right) . \tag{8}
\end{equation*}
$$

We note from (6) that

$$
\begin{equation*}
P_{s_{1} s_{1}}\left(\tau_{n}\right)=\mathrm{e}^{-a\left(1-\sigma_{n}\right)} \sigma_{n}{ }^{s_{1}} L_{s_{1}}\left(-\frac{a\left(1-\sigma_{n}\right)^{2}}{\sigma_{n}}\right) \tag{9}
\end{equation*}
$$

where as before, $\sigma_{n}=z_{1} z_{2} \cdots z_{n}, z_{i}=\mathrm{e}^{-\mu t_{i}}$ for $1 \leqslant i \leqslant n$, and note the fact [6] that

$$
\begin{equation*}
c_{n}(m ; x)=(-x)^{-n} n!L_{n}^{(m-n)}(x) . \tag{10}
\end{equation*}
$$

Using the following well-known generating function for the simple Laguerre polynomials [8]:

$$
\sum_{n=0}^{\infty} t^{n} L_{n}(x)=\frac{1}{1-t} \exp \left(-\frac{x t}{1-t}\right) \quad|t|<1
$$

we obtain from (7) the explicit result

$$
\begin{equation*}
J_{n}=\frac{1}{1-\sigma_{n}} . \tag{11}
\end{equation*}
$$

On the other hand, with (6) substituted for the various transition probabilities in (8), and using the connection between the Charlier and the generalized Laguerre polynomials from (10), we find
$J_{n}=\exp -\left\{a\left(n-\sigma_{1}\right)\right\} \sum_{s_{1}, s_{2}, \ldots, s_{n}=0}^{\infty} Z_{1}^{s_{1}} Z_{2}^{s_{2}} \cdots Z_{n}^{s_{n}} L_{s_{2}}^{\left(s_{1}-s_{2}\right)}\left(X_{1}\right) L_{s_{3}}^{\left(s_{2}-s_{3}\right)}\left(X_{2}\right) \cdots L_{s_{1}}^{\left(s_{n}-s_{1}\right)}\left(X_{n}\right)$
where we have defined

$$
\begin{align*}
& Z_{i}= \begin{cases}\frac{z_{i-1}\left(1-z_{i}\right)}{1-z_{i-1}} & 2 \leqslant i \leqslant n \\
\frac{z_{n}\left(1-z_{1}\right)}{1-z_{n}} & i=1\end{cases}  \tag{13}\\
& X_{i}=-a \frac{\left(1-z_{i}\right)^{2}}{z_{i}} \quad 1 \leqslant i \leqslant n . \tag{14}
\end{align*}
$$

Equating the results in (11) and (12) we then have an expression for a multiple generating function for the product of generalized Laguerre polynomials of the same genre as that given in (1). Note however, that while the forms of the generating function from (11) and (12), and that in (1) are similar, the parameters of $Z_{i}$ and $X_{i}$ in the former case are not independent, their connection being expressed through the intrinsic variable $z_{i}=\exp -\mu t_{i}$ from the relations given in (13) and (14).

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